Controlling A Meandering Wake: Insights From Full-Information Control

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Abstract—In this paper we design and analyze a full-information \(H_\infty\) controller to in order to reduce the wake meandering behind a wind turbine. The low frequency instability that causes wake meandering can unsteady mechanical loads on the downstream turbines resulting in early onset of material fatigue. Controlling the wake meandering in a wind farm can therefore reduce maintenance costs. The control design and analysis in this paper proceeds in two steps. First, a linear reduced order model of the turbine is obtained using snapshots from a higher-order nonlinear 2D actuator-disk model. A full-information \(H_\infty\) controller is then designed for the reduced order model assuming access to the entire flow field and disturbance input. The control performance is evaluated by simulations on the higher-order nonlinear model. The full-information controller can not be implemented in practice. However, it can provide insight into control design for wind farms such as identifying desirable locations to measure the downstream flow.

I. INTRODUCTION

Many states in United States have a regulatory mandate to increase production of energy from renewable sources. Wind energy will be a significant contributor in achieving this target. However, the presence of wake meandering behind turbines in a wind farm can pose a problem in achieving this goal. The meandering wake causes downstream turbines to experience unsteady structural loads which can be damaging and add to the maintenance costs. By controlling this meandering wake therefore, wind energy production can be made more efficient in order to maximize the power of existing wind farms. In this paper we present an ideal, full-information control design to control the wake meandering behind a single turbine. The insights gained from this control design and analysis can be used to design more realistic and implementable controllers for wake control in wind farms.

Wake meandering behind a turbine is characterized as the low-frequency periodic lateral displacement of the wake, generally at downstream distances greater than three times the turbine rotor diameter. It is a well-documented but little understood phenomenon. The exact mechanism causing this meandering instability is not yet known but several theories \([1]\) have been proposed in literature. There is evidence to suggest that the underlying mechanism behind wake meandering is similar to the mechanism behind the instability of helical large-scale coherent structures found behind bluff bodies. For axisymmetric bodies, there is no preferred direction of the meandering. However, the fact that a wind turbine has a specific rotational direction breaks the symmetry which results in a preferred direction of the meandering and at a rather distinct frequency \([2]\). The frequency of meandering appears to depend upon both the thrust coefficient and the tip-speed ratio of the turbine \([3]\).

The literature \([4]\), \([5]\), \([6]\) on wind turbine wake studies shows that, under both steady uniform inflow and more realistic turbulent inflow conditions, approximations of axial turbines as actuator disks or rotating actuator lines successfully produce the wake-meandering instability. Furthermore, there are studies \([7]\), \([8]\) supporting the assumption that the far wake structure predicted using simplified approximations of the turbine geometry provides a good approximation to the wake arising in realistic turbine geometries. Moreover, the results are in reasonable agreement with wind tunnel measurements. The advantage of using these simplified models is that the complete simulation can be run within minutes on a desktop computer and the generated wake is a reasonable approximation to higher-fidelity models.

In this paper, we use the simplifying approximation of the axial turbine as an actuator disk in 2D flow in order to model the wake meandering for our problem (Section II). This nonlinear simulation has approximately 20,000 states and cannot be directly used to design a controller. Therefore we obtain a reduced order linear model using the input output reduced order modeling (IOROM) technique from \([9]\) in Section III. The identified linear IROM preserves the input-output behavior of the nonlinear system and is suitable for control design. We assume access to full state as well as the disturbance input and design a full-information \(H_\infty\) controller for this reduced model in Section IV. The controller performance is validated in full-order nonlinear simulations and the results are presented in Section V. While a full-information controller can not be implemented in an actual wind farm, the analysis nonetheless gives some valuable insights which are also presented in Section V. These insights can be used for sensor placement to give relevant measurements directly to a output measurement based controller. Finally, conclusions and possible directions for future work are given in Section VI.

II. PROBLEM FORMULATION

A. Wind Turbine Setup

Consider a horizontal axis wind turbine with a rotor diameter of 1\(D\) [m] located at an arbitrary location in a rectangular field which spans 20\(D\) in the streamwise \(x\)-direction and 5\(D\) in spanwise \(y\)-direction. The turbine is modeled as an actuator disk with an input axial induction factor \(a\). The axial induction factor is defined as \(a := 1 - \frac{U_{in}}{U}\) with \(u\) denoting the average horizontal air speed across the rotor plane and \(U_{in}\) denoting the average inflow air speed. The power captured from the turbine is given by:

\[
P = \frac{1}{2} \rho U_{in}^2 D^2 a
\]

where \(\rho\) is the air density. The power captured can be expressed as:

\[
P = C_p \frac{1}{2} \rho U_{in}^2 D^2
\]

where \(C_p\) is the power coefficient. The power coefficient can be related to the tip-speed ratio \(\lambda\) and the thrust coefficient \(C_T\) as:

\[
C_p = \left( \frac{C_T + 1}{\lambda} \right)^2
\]

This relationship can be used to control the wind turbine.

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\[ P = \frac{1}{2} \rho A u^3 C_P(a) \]  

where \( \rho \) [kg/m\(^3\)] is the air density, \( A \) [m] is the area swept by the rotor, \( u \) [m/s] is the wind speed perpendicular to the rotor plane, and \( C_P(a) \) is the non-dimensional power coefficient, which is a function of the axial induction factor \( \lambda \):

\[ C_P(a) = 4a(1-a)^2 \]  

As the wind turbine extracts energy from the incoming wind, a wake is generated behind the turbine. The wake interior is characterized by reduced wind speeds and increased turbulence. In the case of a wind farm, we can have a turbine operating in the wakes of upstream turbines and this can result in overall loss of power production and greater structural loading for the downstream turbines. The reduction in velocity immediately downstream of a turbine is directly related to the momentum extracted from the flow which is determined by the turbine thrust coefficient \( C_T(a) \), which is also a function of the axial induction factor \( \lambda \):

\[ C_T(a) = 4a(1-a) \]  

The optimal induction factor that maximizes the power captured from the wind turbine is \( a_0 = \frac{1}{3} \). This optimal induction factor gives rise to a power coefficient of \( C_{P_0} = \frac{16}{27} \) and a thrust coefficient of \( C_{T_0} = \frac{8}{9} \).

On a real turbine, \( C_P \) and \( C_T \) are typically modeled as functions of the tip-speed ratio \( \lambda \) and blade-pitch angle \( \beta \) which can be many-to-one mapping from \((\beta, \lambda)\) to \( C_P \) and \( C_T \) [11]. We use \( C_T(a) \) as an input in this paper which can be mapped to equivalent blade-pitch angle and tip-speed ratio contour. Choosing a specific \((\beta, \lambda)\) pair then depends on other parameters like loads and operating conditions. There is no simple formula for this which is generically applicable to all turbines. However, for a given rotor, a tool like the standalone driver for AeroDyn v15 [12] can be used to compute the \( C_P \) and \( C_T \) as a function of \( \beta \) and \( \lambda \).

### B. Governing Equations

The actuator disk model [13], [14], [15] considered in this paper solves the 2D unsteady, incompressible Navier-Stokes equations assuming a linear drag force acting on the flow due to the turbine. The typical operating wind speeds for a wind turbine do not exceed Mach 0.1 at sea level and hence the assumption of incompressibility is justified. We assume that the freestream flow is orthogonal to the turbine rotor plane and that any disturbances in the freestream act solely to perturb the streamwise flow. We also assume that all velocities are non-dimensionalized by freestream velocity \( U_\infty \) [m/s], all spatial lengths are non-dimensionalized by the turbine diameter \( D \), time \( t \) is non-dimensionalized by \( T = 1 \) s and pressure \( p \) is non-dimensionalized by \( \rho U_\infty^2 \). The dimensionless Navier-Stokes equation governing the evolution of the flow under the assumptions outlined above are given by:

\[
\begin{align*}
\frac{\partial u}{\partial t} + u \frac{\partial u}{\partial x} + v \frac{\partial u}{\partial y} &= - \frac{\partial p}{\partial x} + \frac{1}{Re} \left( \frac{\partial^2 u}{\partial x^2} + \frac{\partial^2 u}{\partial y^2} \right) + f + f_D \\
\frac{\partial v}{\partial t} + u \frac{\partial v}{\partial x} + v \frac{\partial v}{\partial y} &= - \frac{\partial p}{\partial y} + \frac{1}{Re} \left( \frac{\partial^2 v}{\partial x^2} + \frac{\partial^2 v}{\partial y^2} \right) + f
\end{align*}
\]

Here \((u, v)\) denote the streamwise and spanwise velocity components and \((x, y)\) denote the streamwise and spanwise coordinates. The Reynolds number, \( Re \) is defined as \( Re := \frac{U_\infty D}{\nu} \) where \( \nu \) [m²/s] is the kinematic viscosity. \( f \) is the forcing due to the turbine and \( f_D \) is the disturbance forcing acting on the flow.

The forcing due to each turbine is a nonlinear function of space and the time-varying thrust coefficient, i.e. \( f = f(x, y, C_T(t)) \). For modeling purposes, this forcing is split into two components \( f(x, y, C_T(t)) = f_T(x, y, C_{T_0}) + f_C(x, y, C_T(t)) \). Here, \( f_T \) is the constant forcing due to the turbine operating at the a thrust coefficient of \( C_{T_0} = \frac{8}{9} \), which corresponds to optimal axial induction factor of \( a_0 = \frac{1}{3} \). The forcing profile for \( f_T \) is largest at the rotor hub and decreases in magnitude towards the blade tip. The forcing \( f_C \) is due to the variations in the axial induction. This term is modeled as a time-varying input \( C_T(t) \) multiplied by a spatial profile. The input \( C_T(t) \) can be mapped back to an equivalent axial induction factor \( a(t) \). The spatial profile for \( f_C \) is smallest at rotor hub and increases in magnitude outwards towards the blade tip. The disturbance forcing \( f_D \) is a spatial sinusoid at approximately unit distance downstream from the inlet which spans the field in \( y \)-direction and acts to break the symmetry of the incoming flow. Additional details on the forcing functions can be found in the Appendix.

For the simulation of the actuator disk model, the equations are solved using a central differencing scheme [16]. The rectangular region with the turbine is divided into a grid with \( N_x \) points in the \( x \)-direction and \( N_y \) points in the \( y \)-direction. Each grid point is associated with a \( u \) and \( v \) velocity, thus the model has \( 2 \times N_x \times N_y \) states. For our particular example, the field is \( L = 20 \) units long in \( x \)-direction and \( W = 5 \) units wide in the \( y \)-direction. The turbine hub is located at \((x_0, y_0) = (5, 2.5)\). The spacing between the grid points is \( \delta x = 0.1 \) and \( \delta y = 0.1 \) with a time step of \( \delta t = 0.01 \). The resulting model has 20,502 states. The boundary conditions of the model are set to:

\[
\begin{align*}
u(x = 0, y, t) &= u(x, y = 0, t) = u(x, y = W, t) = U_\infty \\
v(x = 0, y, t) &= v(x, y = 0, t) = v(x, y = W, t) = 0
\end{align*}
\]

The Reynolds number for our simulation is set to 50. This small value of \( Re \) is not realistic for wind turbines where the appropriate Reynolds number is of the order of \( 10^6 \). However, for the purposes of this paper, we are restricting ourselves to a low Reynolds number to demonstrate the feasibility and possible insights that can be gained from full-information control of wake meandering.

### C. Wake Meandering

Figure 1 shows an instantaneous pseudocolor plot of the streamwise (top subplot) and spanwise (bottom subplot) flow field in our simulation. For this simulation, the disturbance input is selected to be a uniformly distributed, zero-mean random signal with magnitude 0.3. The excitation input, \( C_T \) is a sum of sines at 1.1 and 1.25 rad/s superimposed with a uniformly distributed, zero-mean random noise of magnitude 0.1. The wake meandering as evidenced by the characteristic periodic lateral displacement of the wake behind the turbine can be clearly seen from the plot.
This meandering wake can cause many problems in a wind farm. Downstream turbines located in the path of such oscillating wakes experience fluctuating inflow conditions. The changing inflow conditions adversely affect the structural loads on downstream turbines. The unsteady loads can cause early onset of material and structural fatigue thereby adding to maintenance costs. Wake control can therefore give significant revenue boost by reducing the costs of operation. In the following sections, we obtain a reduced order model from the full-order actuator disk simulation and design a controller for the turbine wake meandering.

III. MODEL REDUCTION

The actuator disk model obtained in the Section II-B has more than 20,000 states, and as such is not suitable for control design. The nonlinearity of the model adds another complication as the extensive theory available for control of linear models can not be applied [17]. Therefore, we first construct a reduced order linear model that captures the dominant input-output behavior of the nonlinear model. There are several techniques available in literature for reducing the model order, such as balanced truncation [18], [19], proper orthogonal decomposition (POD) [20], [21], balanced POD [22], [23], [24], dynamic mode decomposition [25], [26], [27], and input-output reduced order models (IOROMs) [28], [29], [9]. A good overview of existing approaches can be found in [30].

For our problem we use the model reduction approach from [9] to obtain a two-input one-output linear model. The two inputs are the freestream disturbance and turbine thrust coefficient. The disturbance input is included in order to model its effect on the wake meandering. The thrust coefficient acts as the control input. The single output is the measurement of spanwise velocity \( v \) at \((x_M, y_M) = (13, 2.5)\). This output is convenient to observe the fluctuations in \( v \) due to wake meandering. The model is constructed at a single operating point defined by the non-dimensional freestream velocity \( U_{\infty} \) and optimal turbine input. In order to obtain the linear reduced order model, we excite the nonlinear system about the chosen operating point using the inputs and gather state, input and output snapshots from the simulation. One snapshot is gathered for every 20 steps of the nonlinear simulation giving an effective time step of \( 20f = 0.2 \). A lower-order projection basis for the states is obtained using proper orthogonal decomposition and then a linear model fit is obtained for the reduced state, input and output snapshots. The reduced order system matrices are given by:

\[
\begin{bmatrix}
F & G \\
H & D
\end{bmatrix} = 
\begin{bmatrix}
Q^T X_1 \\
Q^T X_0
\end{bmatrix}
\tag{4}
\]

where \( F, G, H, D \) are the matrices for the reduced order model, \( Q \) is the set of POD projection modes, and \( X_0, X_1 \), \( Y_0 \) and \( Y_1 \) are respectively the input and output snapshots. \( \dagger \) denotes the Moore-Penrose pseudoinverse of the matrix. Further details of the algorithm can be found in [9]. The approach can be extended to capture wider operating conditions using parameter-varying models [17] and gain-scheduled control.

![Figure 1. Flow-field of the 20, 002 state simulation at \( t = 45.0 \). First subplot shows \( u \) and second subplot shows \( v \). The vertical black line represents the location of the actuator disk.](image1.png)

![Figure 2. Bode plots of the identified reduced order linear model.](image2.png)

A 34 state reduced order model is identified using this approach. The Bode plot of the identified reduced order model is given in Figure 2. In the figure, \( G \) corresponds to transfer function from \( C_T \) and \( G_d \) corresponds to transfer function from freestream disturbance to the measured output. The time step of 0.2 between snapshots gives a Nyquist frequency of \( \omega_h = 15.71 \) rad/s which is approximately 10 times the observed wake meandering frequency (\( \approx 1 – 1.5 \) rad/s). The total time of simulation was chosen to be 100 which sets the minimum identifiable frequency of \( \omega_l = 0.06 \) rad/s. The identification signal for input \( C_T \) was selected to be a chirp signal with frequency range \( [\omega_h/2, \omega_h/2] \) rad/s and magnitude 1/9 while a uniform random noise with magnitude 0.3 was selected as the disturbance input. Thus the identified model is valid between approximately 0.06 rad/s and 7.85 rad/s.

The validation results of the IOROM model for the same input as that described in Section II-C are given in Figure 3. The first subplot shows the output from the nonlinear system vs. the output obtained from IOROM. The green vertical line represents the time instant \( t = 45.0 \) for which the subsequent pseudocolor subplots are drawn. The pseudocolor plots in Figure 3 are obtained by lifting up the reduced order states to obtain approximations for the full-order nonlinear states. The reduced order model is able to adequately capture the nonlinear output as well as behavior of full-order system states as can be clearly seen by comparing the pseudocolor plots in Figure 3 to Figure 1.

IV. FULL-INFORMATION \( H_\infty \) CONTROLLER

The next step is to synthesize a controller using the reduced order model. Our initial designs could not reduce the wake...
meandering using the single measurement at \((x_M, y_M)\). This section focuses on a full-information design to gain insight into additional measurements that are most beneficial. In particular, we investigate a full-information \(H_\infty\) controller that has access to all the states as well as the disturbance.

The matrix information required to reduce the effect of the disturbance on the measurement vector includes the full state as well as are constant system matrices of appropriate dimensions. Note found by solving the following discrete-time algebraic Ricatti equation:

\[
\begin{align*}
Y_k & = (D^T_1 D_1 + B^T_1 P B) (B^T_1 P A + D^T_1 C), \\
F_2 & = -(D^T_1 D_1 + B^T_1 P B)^{-1} (B^T_1 P B D_2 + D^T_1 D_2).
\end{align*}
\]

Additional details can be found in [31]. This result is applied using the reduced order model for the actuator disk dynamics and hence \(n_x = 27\). There is a single control input at the turbine \((n_u = 1)\) and a single disturbance near the upstream boundary condition \((n_d = 1)\). Finally, the error, i.e., the signal to be minimized, is also scalar \((n_e = 1)\) and given by the downstream lateral velocity. All these signals are as described in Section III. The next section discusses the performance of this full-information controller. In addition, the controller matrices \(F_1\) and \(F_2\) are used to gain insights that can be used to identify the information that is most important for wake control.

V. RESULTS AND DISCUSSION

A full-information controller was designed using an augmented plant \(P\) obtained by augmenting the reduced order model with a weight \(W = \frac{64}{27} (\pi + 0.005)^2\) to penalize the control input at high frequencies.

\[
P = \begin{bmatrix} W & 0 \\ G & G_d \end{bmatrix}
\]

Note that the implementation of controller on the nonlinear model requires a projection of the full-order state down to the reduced order state using the projection modes. For the plant \(P\), the controller can be written as \(u_k = F_1 Q T x_k^{\text{full}} + F_2 d_k + F_3 x_k W\), where \(x_k^{\text{full}}\) is the full-order system state, \(x_k^W\) is the state of weight \(W\), and \(Q^T\) contains the projection modes.

The full-order, nonlinear, actuator disk model was simulated with the controller and the results are shown in Figure 4. For this simulation, the actuator disk model was initialized at the base flow conditions and the input disturbance was designed to be a uniformly distributed, zero mean random noise signal of magnitude 0.3. The first subplot shows the control input \(C_T\) relative to the trim value of \(C_T_0 = \frac{2}{3}\). As the total thrust coefficient cannot exceed 1, the limiting positive value of control input is \(1 - C_T_0 \approx 0.11\). The required control effort never exceeds 0.06 so clearly the controller does not demand
excessive control authority. The second subplot compares the output of the system without control (red line) with the controlled output (blue dashed line). It can be seen that the controller successfully suppresses the wake meandering.

The ideal, full-information controller obtained in Section IV can intuitively be understood to give the best possible control performance as there is no information hidden from the controller. While such a controller is not realistically implementable, the exercise provides several valuable insights. The first term $F_1 Q^T x_{\text{full}}^T$ of the control input can be understood as $K x_{\text{full}}^T$ where $K = F_1 Q^T$ is the control gain applied to full-order states. The distribution of values in $K$ therefore informs about the states which are most important for control purposes. In other words, the states for which the control gain is high are most important to the full-information controller.

Figure 5 shows the pseudocolor plots of $K$ for both streamwise and spanwise velocity components. There appear to be four locations which are important for $u$ and two locations which are most important for $v$. These points of importance, i.e. large magnitude gains, correspond to the areas of sharp red or blue color downstream of the turbine. We think that the four highlighted gains in the streamwise direction might be the controller attempting to act on the vorticity of the flow. However, further investigations are underway.

![Fig. 5. Controller Insight from distribution of control gain $K$](image)

We hypothesize that using only six measurements at locations corresponding to high gain as inputs for an output measurement based controller we might be able to recover most of the full-information control performance. Another approach would be to try to estimate these states from measurements at other, perhaps more feasible locations and use those estimates for an observer-based control design.

To test our hypothesis, we made the disturbance gain $F_2 = 0$ to neglect the disturbance feedback and then gradually zeroed out the entries in the gain matrix $K$ starting from the smallest to the largest in magnitude. As $K$ is gradually emptied, the gains at those six locations are preserved the longest (being largest in magnitude) while the smaller gains are made zero. Using these gain matrices, the simulation was rerun and the variance in output was calculated. Figure 6 shows the semi-log plot of variance of controlled output (normalized with respect to the open loop variance) vs. the log of percentage of non-zero entries in $K$. The results are encouraging and it can be seen that with only $\approx 15\%$ non-zero entries in $K$, the variance in output is better than that obtained using the full gain matrix. When $\approx 100\%$ of the entries of $K$ are zero, the system is essentially running with no control and the normalized closed loop variance approaches 1. More formal sparsity-promoting techniques [32] will be investigated in the future for obtaining a sparse $K$.

![Fig. 6. Variance in measured output (normalized w.r.t. open-loop variance) as the entries of gain matrix $K$ are zeroed out](image)

VI. CONCLUSIONS

This paper considers the control of wake meandering behind a turbine modeled with a simplified, nonlinear 2D actuator disk model. A linear reduced order model is constructed from input-output data to approximate the full-order nonlinear simulation. This reduced order model is then used to design a full-information controller which is successfully validated in the nonlinear simulation. The control gains of the full-information controller are used to identify the information most useful for control. This approach can also be used for analyzing a full-information controller for a wind farm model in Simulator fOr Wind Farm Applications (SOWFA). The insights gained from high-fidelity models like SOWFA can be valuable in deciding measurement locations and designing controllers for actual wind farms, which also forms a direction for further future work.

APPENDIX

A. Turbine Forcing

As the hub of the turbine is placed at $(x_0, y_0)$, the rotor plane lies within $(y_0 - \frac{1}{2}) \leq y \leq (y_0 + \frac{1}{2})$. The forcing terms introduced by the turbines are defined as:

$$f_T(x, y, C_{T_0}) = 0.7 C_{T_0} C_{\theta_v} C_{\theta_v} (1 - |\Delta_x|)(1 - |\Delta_y|)^{0.7}$$

$$f_C(x, y, C_T(t)) = C_T(t) \text{sign}(\Delta_y)(1 - |\Delta_x|)|\Delta_y|^{0.7}$$

for $(x_0 - 2\delta x) \leq x \leq (x_0 + 2\delta x)$ and $(y_0 - \frac{1}{2} - 2\delta y) \leq y \leq (y_0 + \frac{1}{2} + 2\delta y)$.

For all other values of $x$ and $y$, $f_T = 0$ and $f_C = 0$. Here $\Delta_x = (x - x_0)$ is the $x$-direction displacement from the hub center, $\delta x$ is the $x$-spacing between grid-points, and

$\delta y$ is the $y$-spacing between grid-points, and

1It is important to note that, in general, modifying the control gain matrix can cause the closed loop system to go unstable.

2This increase in performance is because we are optimizing with respect to the $H_\infty$ norm but then doing the performance evaluation (Figure 6) with respect to output variance.
\[ C_{T_0} = \cos \frac{3\Delta x}{25\pi} \] works to smooth the transition from unforced to forced region in the flow field. \( \Delta y \), \( \delta y \) and \( C_{T_0} \) are defined analogously in terms of \( y \). \( C_{T_0} \) is the thrust coefficient of the turbine operating the optimum axial induction factor of \( a = \frac{1}{4} \), and \( C_T(f) \) is the variation of thrust coefficient about this optimum value.

### B. Disturbance Forcing

The disturbance forcing is intended to break the spatial symmetry of the free stream. The forcing acts at approximately 1D downstream of the flow field inlet. The disturbance forcing can be computed as per the below pseudocode:

```plaintext
Initialize \( f_D = \text{zeros}(N_x, N_y) \)
for \( i = 1 \) to \( N_y = 4 \)
for \( j = 1 \) to \( N_x = 9 \)
for \( k = 1 \) to \( 4 \)
\( x = k + 9 \) \( y = N_y \times (i + j - 1) \)
\( C_x = 0.25 \times (1 + \cos ((k - 2.5) \frac{\pi}{y})) \)
\( C_y = 0.25 \times (1 + \cos ((j - 2.5) \frac{\pi}{y})) \)
\( f_D(x, y) = f_D(x, y) + g(y)C_xC_y \)
end loop of \( k \)
end loop of \( j \)
end loop of \( i \)
```

where \( g \) is given by:

\[
g(y) = \sum_{i=1}^{8} \sin(2\pi \left( \frac{y}{N_y} + \text{rand} \right))
\]  

(11)

Here \( \text{rand} \) is a random number between 0 and 1.

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